

# Optics and Laser Technology

## Controlling the shell thickness of SiO<sub>2</sub> on TiO<sub>2</sub> NPs: Characterization, linear and nonlinear optical properties

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## Highlights

1. Core-shell  $\text{TiO}_2@\text{SiO}_2$  nanoparticles with different shell thicknesses were synthesized using the Stöber method.
2. The impact of the silica shell thickness on the structural and optical properties of  $\text{TiO}_2$  was investigated.
3. Linear and non-linear optical parameters of  $\text{TiO}_2$  and  $\text{TiO}_2@\text{SiO}_2$  with different shell thickness were investigated.

## Controlling the shell thickness of SiO<sub>2</sub> on TiO<sub>2</sub> NPs: Characterization, linear and nonlinear optical properties

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### Abstract

TiO<sub>2</sub> nanoparticles (NPs) were prepared by the sol-gel method and then encapsulated with different thicknesses of SiO<sub>2</sub> as a shell using the Stöber method. The crystal, chemical structure, and morphology of the core-shell TiO<sub>2</sub>@SiO<sub>2</sub> NPs were characterized using XRD, FTIR, and TEM techniques. The amorphous nature of the shell (SiO<sub>2</sub>) was shown by XRD examination, which influenced the crystallinity of the TiO<sub>2</sub> NPs, while FTIR data verified the association between SiO<sub>2</sub> and TiO<sub>2</sub>, and a TEM study validated the coating of TiO<sub>2</sub> NPs with SiO<sub>2</sub>. The optical absorption of TiO<sub>2</sub> NPs was characterized by a sharp absorption edge around 343 nm, which redshifted to a higher wavelength region with increasing shell thickness. The direct bandgap was found to decrease from 3.97 eV for TiO<sub>2</sub> to about 3.60, 3.37, and 3.26 eV for TiO<sub>2</sub>@SiO<sub>2</sub>, with shell thicknesses of about 2.5, 5.5, and 8 nm, respectively. The refractive index and extinction coefficient increased with increasing shell thickness. Furthermore, an enhancement in the optical conductivity of TiO<sub>2</sub> NPs was observed with increasing shell thickness. The Wemple-DiDomenico model was applied to compute the nonlinear refractive index ( $n_2$ ), first- and third-order ( $\chi^{(1)}$  &  $\chi^{(3)}$ ) susceptibilities values. An enhancement was observed for all nonlinear optical parameters with the insertion of the shell SiO<sub>2</sub> into the core TiO<sub>2</sub>. The value of ( $\chi^{(3)}$ ) was increased from  $2.08 \times 10^{-14}$  for pure TiO<sub>2</sub> NPs to  $3.04 \times 10^{-09}$  for the core-shell TiO<sub>2</sub>@SiO<sub>2</sub> sample with a shell thickness of 8 nm.

**Keywords:** Nanoparticles; Sol-gel; TiO<sub>2</sub>; TiO<sub>2</sub>@SiO<sub>2</sub> core/shell; Refractive index; Wemple–DiDomenico model.

## 1. Introduction

Recently, there has been growing interest in the development of nanomaterials for biomedical applications due to their unique properties and potential for enhancing diagnostic and therapeutic approaches [1-6]. Among these nanomaterials, core-shell nanoparticles (NPs) have emerged as a promising platform for various biomedical applications [7-10]. In particular, core-shell  $\text{TiO}_2@\text{SiO}_2$  NPs have attracted significant attention due to their combined properties and versatile applications in the field of biomedicine [11-14]. The core of these NPs is composed of metal oxide nanocrystalline semiconductive materials, particularly titanium dioxide ( $\text{TiO}_2$ ), known for its excellent photocatalytic and optical properties [14-18].  $\text{TiO}_2$  NPs exhibit high photostability [19,20] and have attracted significant attention across various fields due to their vast potential in photocatalysis [21-24], solar cells [25-27], nonlinear optics [28-31] and sensors [32-35]. Due to the chemical and electrostatic stability of silica, it can operate as a protective shell to limit contact with the surrounding media after coating the nanoparticles [36]. Lately,  $\text{SiO}_2$  has been coated with a shell from  $\text{TiO}_2$  and found to enhance the scattering efficiency in dye-sensitized solar cell applications and show enhanced power conversion efficiency [37]. Furthermore, there are many limitations to use bare  $\text{TiO}_2$  nanoparticles in biological systems due to their potential cytotoxicity and lack of stability in physiological conditions [38-42]. To overcome these limitations, a silica ( $\text{SiO}_2$ ) shell is applied as a protective coating around the  $\text{TiO}_2$  core. The  $\text{SiO}_2$  shell provides biocompatibility, stability, and surface functionalization capabilities to the nanoparticles. The biocompatible nature of  $\text{SiO}_2$  allows for the safe use of core-shell  $\text{TiO}_2@\text{SiO}_2$  in biological environments, reducing the risk of adverse effects on cells and tissues [40, 43-45]. Moreover, the  $\text{SiO}_2$  shell offers opportunities for surface modification with various functional groups, biomolecules, or targeting ligands. This surface functionalization enables specific targeting and interaction with biological entities, such as cells, tissues, or biomarkers, making core-shell  $\text{TiO}_2@\text{SiO}_2$  NPs versatile tools for targeted drug delivery, bio-imaging, and sensing applications [45-49]. Furthermore, the unique optical properties of core-shell  $\text{TiO}_2@\text{SiO}_2$  NPs, including light absorption and emission, make them excellent candidates for photodynamic therapy with improvement of biocompatibility [50,51]. They can be utilized as contrast agents in various imaging modalities, including fluorescence imaging, photoacoustic imaging, and computed tomography (CT), allowing for real-time monitoring and visualization of biological processes [52-54].

Core-shell  $\text{TiO}_2@\text{SiO}_2$  synthesis poses a significant challenge, with preparation methods categorized into liquid-solid [46, 55] and gas-solid transformations [46, 56]. Among these methods, the sol-gel technique stands out as a widely used and cost-effective approach for material preparation. It offers several advantages, including easier control over chemical compositions, morphology, and crystallite size of the nanomaterial [14, 46]. Additionally, the sol-gel technique ensures compositional homogeneity, enabling precise regulation of structural morphology and grain size [40, 57, 58].

Herein, a straightforward technique for synthesizing  $\text{TiO}_2@\text{SiO}_2$  core-shell was presented with control of silica shell thickness. Three different thicknesses of the shell were obtained and its impact on the structure and optical properties of  $\text{TiO}_2$  NPs was investigated. Changes in structure and morphology were assessed using XRD, FTIR, and TEM techniques. With the use of UV-visible data, the optical properties of the  $\text{TiO}_2@\text{SiO}_2$  nanocomposite were evaluated.

## **2. Experimental and procedure**

### **2.1. Materials and chemicals**

Titanium (IV) isopropoxide (TTIP) and Tetraethoxysilane (TEOS,  $\geq 99.0\%$ ) were purchased from Sigma-Aldrich, Germany. Isopropanol (IPA) was obtained from Chem-Lab Analytical based in Belgium. Merck in Darmstadt, Germany provided ammonia aqueous solution (25%), Absolute ethanol (99.5%), and nitric acid (69%). All compounds were of reagent quality and were utilized as supplied, with no extra purification required.

### **2.2. $\text{TiO}_2$ Nanoparticles preparation**

$\text{TiO}_2$  NPs were produced using sol-gel technique according to the procedure previously mentioned [30]. Titanium (IV) isopropoxide (TTIP) was dissolved in isopropanol, resulting in the rapid formation of a transparent stock solution. The stock solution was carefully added in small droplets to deionized water while vigorously stirring at room temperature, leading to the formation of a white gel. The solution was subsequently treated with a 1M nitric acid solution until the pH value reached 2 and then maintained at a temperature of  $70^\circ\text{C}$  for duration of 24 hours.

### **2.3. Synthesis of $\text{TiO}_2\text{-SiO}_2$ core-shell nanoparticles**

The Stöber method was used to synthesize  $\text{TiO}_2@\text{SiO}_2$  core-shell [38]. The specific procedure involved several steps. Initially, a mixture of 10 ml deionized water and 100 ml

ethanol was prepared, to which 1 g of  $\text{TiO}_2$  was added. Then, for 30 minutes, the mix was subjected to ultrasonics to disperse the  $\text{TiO}_2$  particles. In the subsequent stage, different concentrations of TEOS (Tetraethyl orthosilicate) were sequentially introduced into the suspension. The TEOS concentrations used were 0.099M, 0.19M, and 0.28M. Additionally, ammonia was added gradually, aiming to attain the desired pH level of 10. This addition of ammonia was carried out under continuous stirring. The resulting mixture was vigorously stirred at 20°C for 12 hours to allow for a complete reaction. Following a 12-hour stirring period, the mixtures underwent centrifugation at 7500 rpm/5 minutes. The liquid phase was carefully removed, leaving behind the resulting silica coated  $\text{TiO}_2$  NPs. To eliminate any excess reactants, the nanoparticles were subjected to three washes with ethanol. Subsequently, the particles were dried overnight at 100 °C. The denotation given to the resulting silica coated  $\text{TiO}_2$  nanoparticles was TSM, where M represents different concentrations of TEOS. Specifically, M was assigned values of 1, 2, and 3, corresponding to TEOS concentrations of 0.099M, 0.19M, and 0.28M (TS1, TS2, and TS3), respectively. The reaction process is visually depicted in Fig. 1.

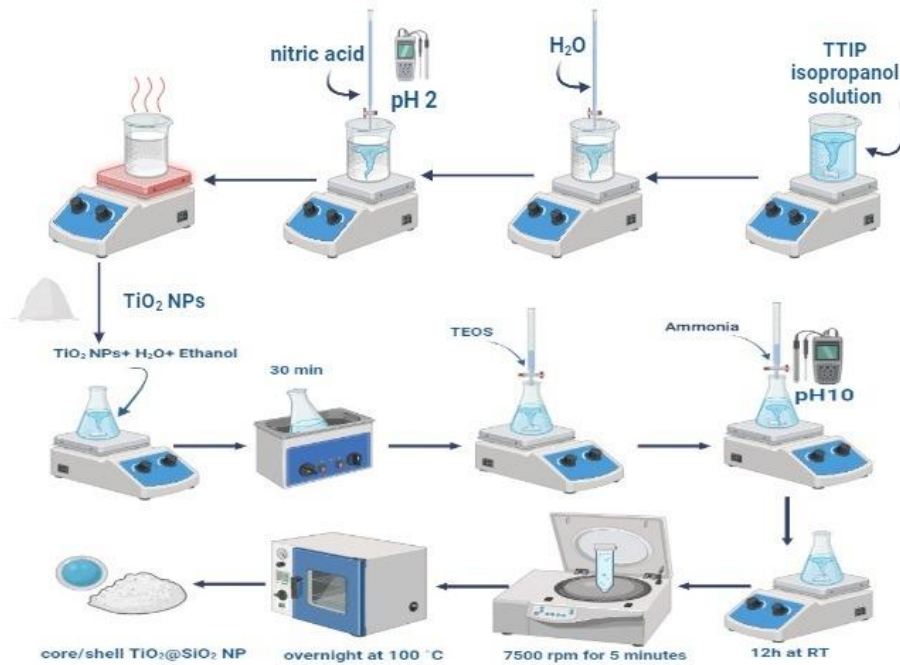


Fig. 1: visual representation of the synthetic pathway utilized to generate core-shell  $\text{TiO}_2@\text{SiO}_2$  NPs.

## 2.4. Material Characterization

The X-ray powder diffraction (XRD) technique was performed using the XPERT-PRO to record the crystal structures of the prepared samples. The XRD measurements were conducted

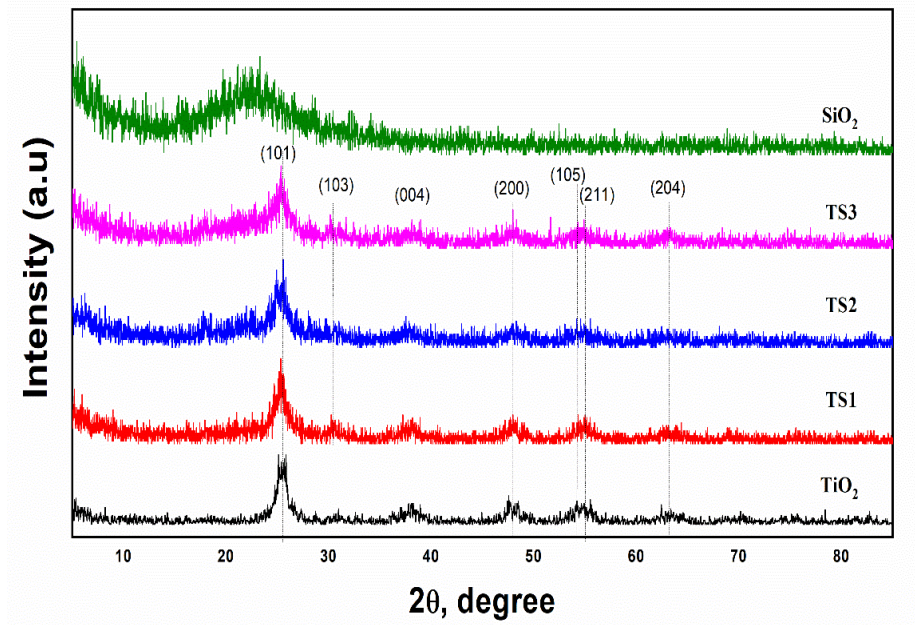
over a range of  $5^\circ \leq 2\theta \leq 90^\circ$ , with a minimum step size of  $2\theta$ : 0.001, and at wavelength ( $K\alpha$ ) =  $1.54614^\circ$ . The morphological characteristics of the produced  $\text{TiO}_2$  powder and synthesized  $\text{TiO}_2@\text{SiO}_2$  core-shell powder was examined using transmission electron microscopy (TEM) with Talos F-200i (Thermo-Scientific) high-resolution TEM operating at an accelerating voltage of 200 kV. To investigate the various chemical bonds presented in the prepared samples, FT-IR vertex 70 RAM II, Bruker was utilized. FTIR spectra were collected at room temperature, covering a  $400\text{-}4000\text{ cm}^{-1}$  range. A double-beam spectrophotometer (Perkin Elmer Lambda 40) was used to analyze the UV-visible absorption spectra of materials.

### 3. Results and discussions

#### 3.1. Characterization

##### 3.1.1. XRD analysis

XRD analysis was performed to validate the crystal structure of the prepared nanoparticles. Fig. 2 presents the XRD patterns of  $\text{TiO}_2$  and  $\text{TiO}_2@\text{SiO}_2$  NPs.



**Fig.2: XRD patterns for prepared samples.**

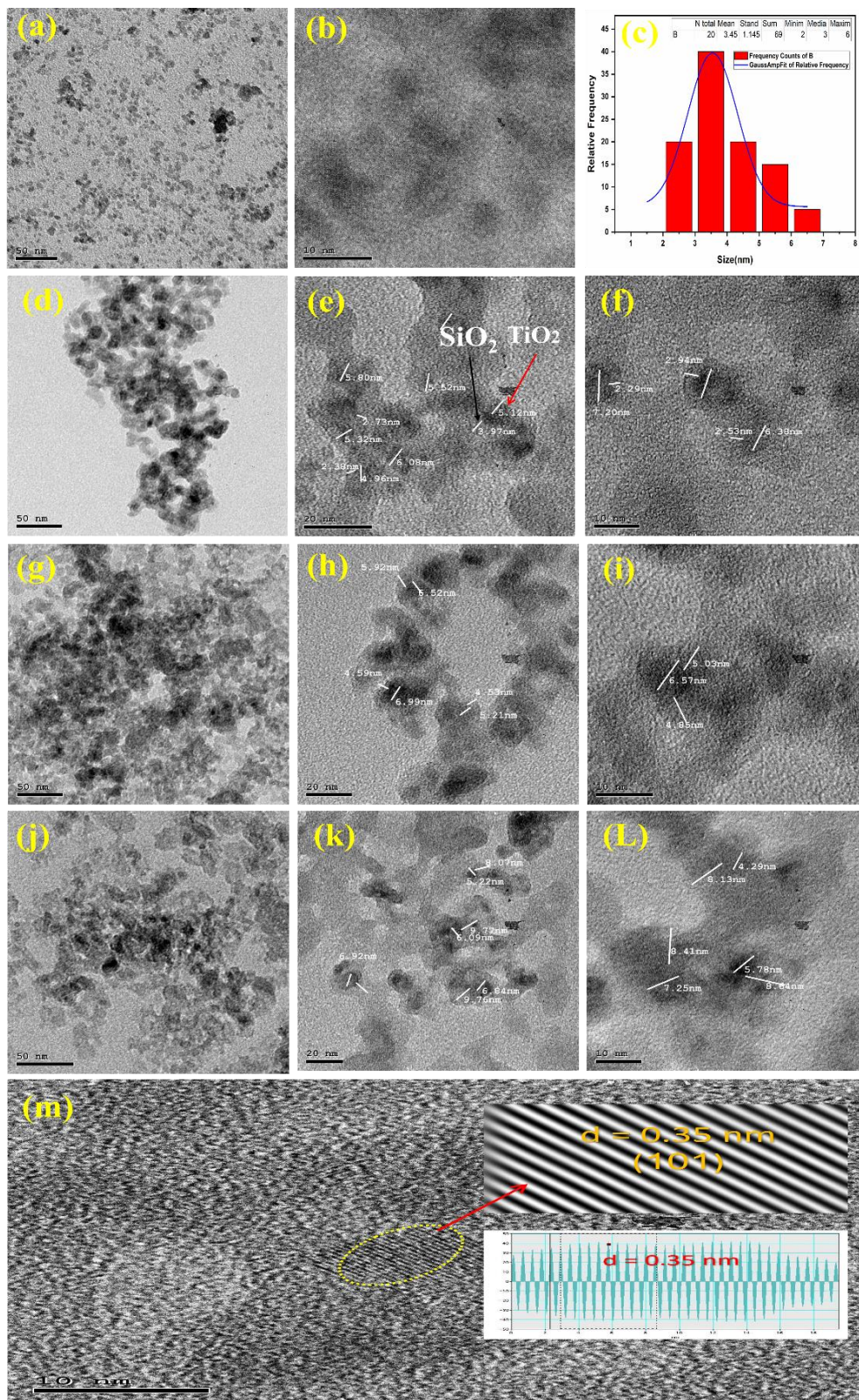
The  $\text{TiO}_2$  pattern coincides with anatase  $\text{TiO}_2$  (JCPDS No. 01-075-2552), with no diffraction peaks correlated to the  $\text{TiO}_2$ -polymorphs identified. Specifically, the peaks at  $25.335^\circ$ ,  $36.98^\circ$ ,  $37.809^\circ$ ,  $48.104^\circ$ ,  $53.92^\circ$ ,  $55.138^\circ$  and  $62.752^\circ$  corresponded to the (101), (103), (004), (200),



(105) and (211) and (204) planes of anatase  $\text{TiO}_2$ , respectively [59]. For pure silica NPs, a broad peak was observed around  $22^\circ$  was assigned to the silica amorphous structure. Diffraction patterns obtained from  $\text{TiO}_2@\text{SiO}_2$  core-shell prepared with different TEOS concentrations further demonstrated the influence of TEOS on the crystallinity of the composite. In these patterns, the diffraction background of TS1, TS2, and TS3 wasn't a flat-line at  $2\theta = 25.335^\circ$  but instead displayed broadband. This was attributed to the presence of amorphous  $\text{SiO}_2$  on the surface of the nanocrystalline  $\text{TiO}_2$  [60, 61]. Moreover, as the TEOS concentration increases, the peak intensity in TS1, TS2, and TS3 becomes weaker and broader, indicating lessened crystallinity [60]. Furthermore, the increase in TEOS concentration promotes the formation of a thicker and more continuous  $\text{SiO}_2$  shell around the  $\text{TiO}_2$  core and induces the formation of amorphous regions within the composite, which results in a broad peak in the XRD spectra [62].

### 3.1.2. Transmission electron microscope analysis

HR-TEM was used to identify the morphology and size of the core  $\text{TiO}_2$  and the shell thickness  $\text{SiO}_2$ , as shown in Fig. 3. The TEM images, specifically Fig. 3a & b, confirmed that  $\text{TiO}_2$  NPs exhibited a spheroidal shape with an average particle size of 3.45 nm, as illustrated by the corresponding histogram (Fig. 3c). For TS1, TS2 and TS3 a clear contrast was observed suggesting that formation of a shell layer from amorphous  $\text{SiO}_2$  around  $\text{TiO}_2$  NPs. Briefly, TEM images of TS1 were displayed in Fig. 3(d, e, and f) having an average size of around 5:6.4 nm with shell thickness  $\sim 2.5$  nm. Whereas TEM images for TS2 shown in Fig. 3(g, h, and i) presented an average size ranging from 9:12 nm with a shell layer of about  $\sim 5.5$  nm. Furthermore, TEM images of TS3 are shown in Fig. 3(j, k, and l). The coated particles displayed a size range of 12:15 nm, with a  $\text{SiO}_2$  thickness around  $\sim 8$  nm. The HR-TEM of TS3, Fig. 3m, showed d-spacing of 0.35 nm that corresponding to (101) plane of the  $\text{TiO}_2$ -anatase.



**Fig.3: TEM images for; (a, b, & c) TiO<sub>2</sub>, (d, e, & f) TS1, (g, h, & i) TS2, (j, k, & l) TS3, and (m) HR-TEM for TS3.**

### 3.1.3. FT-IR analysis

FTIR spectra were conducted on the SiO<sub>2</sub>, TiO<sub>2</sub>, TS1, TS2, and TS3 samples, as depicted in Fig. 4., pure SiO<sub>2</sub> displayed strong absorption peaks at approximately 1069.15 and 978.8 cm<sup>-1</sup> corresponding to the symmetric Si-O-Si bending vibration. Additionally, the absorption bands at 804.57 and 451.95 cm<sup>-1</sup> are assigned to Si-O-Si asymmetric stretching vibrations[68]. For pure TiO<sub>2</sub>, the bands in the region from 400 to 800 cm<sup>-1</sup> are assigned to Ti-O vibration[69], and the de-formative vibration of Ti-OH stretching mode was observed at 1631.70 cm<sup>-1</sup> [70]. For composite structures, a weak absorption band at 954.42 cm<sup>-1</sup> in TS1 is attributed to Si-O-Ti flex vibrations, which testifies to the chemical bond between TiO<sub>2</sub> and SiO<sub>2</sub> [63, 64]. Furthermore, this band was shifted to 956.13 and 957.31 cm<sup>-1</sup> for TS2 and TS3, respectively, indicating the variation in shell thickness for the three composites. FT-IR bands revealed that these absorption peaks were generally more pronounced in the coated TiO<sub>2</sub> samples with higher silica content. This strongly suggests the encapsulation of TiO<sub>2</sub> NPs within the silica matrix [65, 66]. Furthermore, absorption bands in the range of 800–400 cm<sup>-1</sup>, corresponding to the characteristic vibrational modes of Ti-O of TiO<sub>2</sub>, were observed, indicating the presence of the core-TiO<sub>2</sub> structure [38, 67]. Such a result is predicted to have a noticeable impact on the optical properties of the TiO<sub>2</sub> NPs.

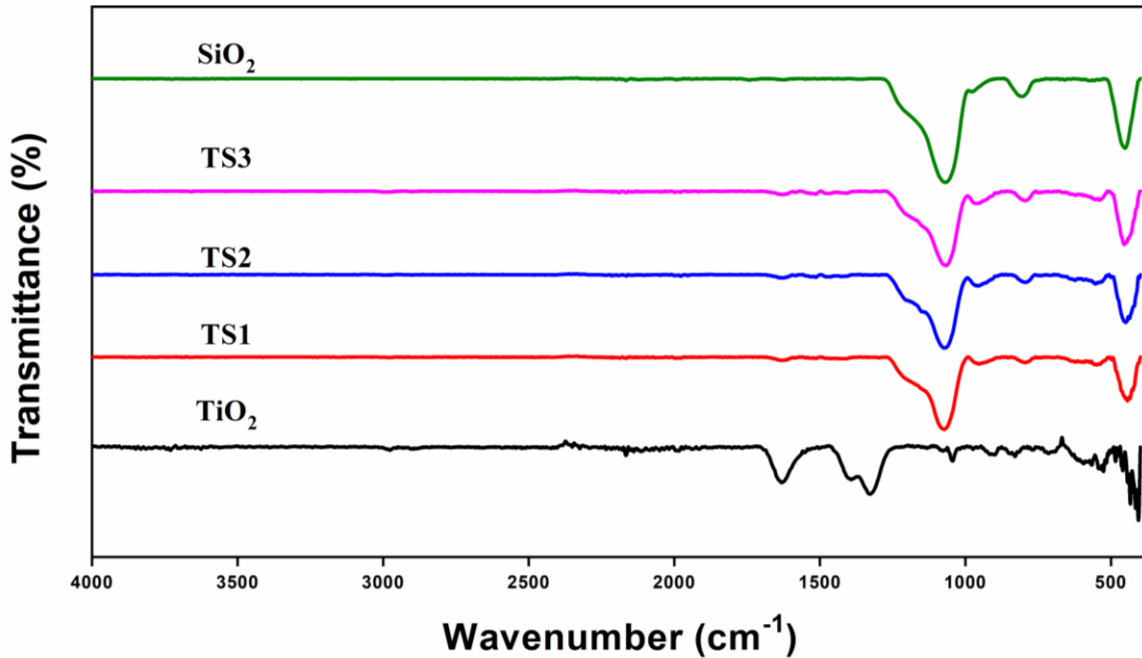


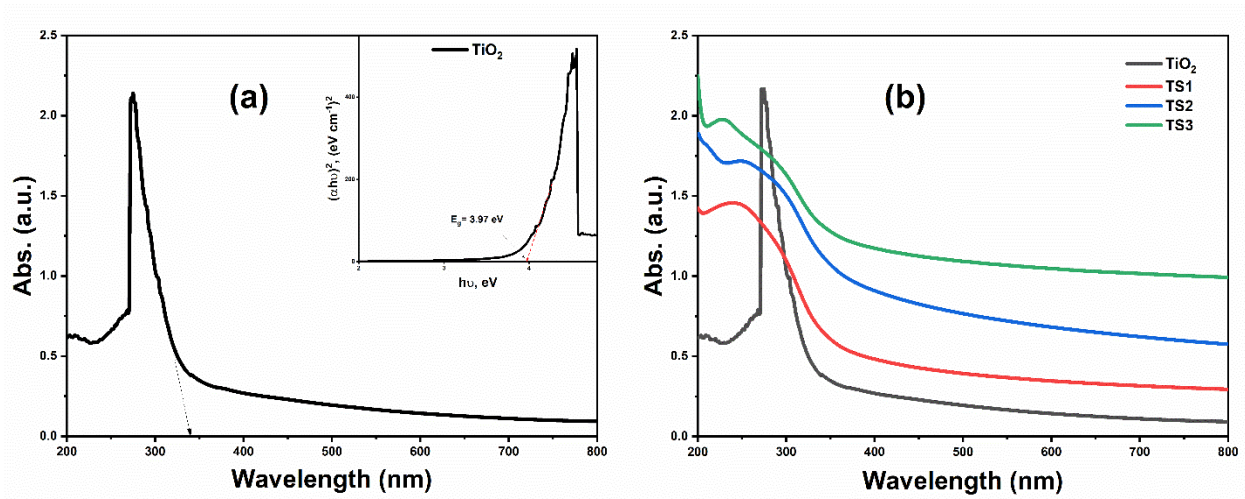
Fig.4: FTIR patterns for prepared samples.



## 3.2 Optical features

### 3.2.1 UV-visible spectra, optical bandgap, and linear refractive index

The TiO<sub>2</sub>@SiO<sub>2</sub> core-shell absorption spectra with different thickness of the SiO<sub>2</sub> shell are shown in Fig.5. The absorption spectrum of the pure TiO<sub>2</sub> NPs shows an absorption edge around a 343 nm range (Fig.5a). Encapsulation of the TiO<sub>2</sub> surface by the silica shell leads to increased absorption intensity and redshift in absorption edge to about 387, 445, and 482 nm for TiO<sub>2</sub>@SiO<sub>2</sub> with shell thickness 3, 6, and 9 nm, respectively (Fig.5b). In addition, the shift increases to a higher wavelength region with increasing the thickness of the SiO<sub>2</sub> shell as shown in Fig.5b. This finding increases the efficiency with which simulated sunlight is used in solar cell and photocatalysis applications.



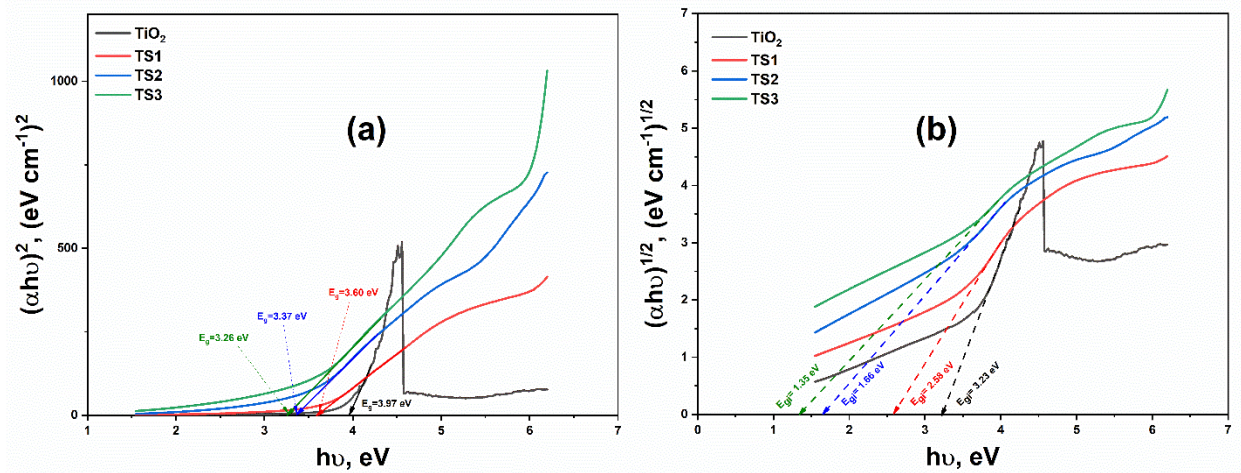
**Fig.5: The absorbance as a function of wavelength for prepared samples.**

The redshift in the absorption spectra is caused because of the quantum confinement effect. As shown in the TEM analysis the SiO<sub>2</sub> shell increases in size with increasing the TEOS concentration during the reaction process. This increase in particle size causes a redshift in the absorption spectra, and the electronic transition structure of the material changes in accordance. This change was discussed with the help of Tauc's relation as follows [71];

$$(\alpha h\nu)^2 = A(h\nu - E_{gd}) \quad (1)$$

$$(\alpha h\nu)^{1/2} = A(h\nu - E_{gi}) \quad (2)$$

where  $E_{gd}$  and  $E_{gi}$  are the direct and indirect bandgap energies, respectively. Fig.6 shows the relation of  $(\alpha h\nu)^2$  and  $(\alpha h\nu)^{1/2}$  as a function of  $(h\nu)$ . The  $E_{gd}$  and  $E_{gi}$  values are extracted in Table 1.



**Fig.6: Plots of (a)  $(\alpha h\nu)^2$  and (b)  $(\alpha h\nu)^{1/2}$  vs.  $(h\nu)$  of the prepared samples.**

as shown in Table 1,  $E_{gd}$  and  $E_{gi}$  values decreased dramatically with increasing the shell thickness. Herein, the  $E_{gd}$  value of the pure  $\text{TiO}_2$  is higher than that of the bulk  $\text{TiO}_2$  (3.2 eV), which could be caused by the tiny size of the  $\text{TiO}_2$  NPs (3 nm) compared to that present in the literature [72-75]. The particle size of  $\text{TiO}_2$  is about 3 nm, and the shell varies from 3 to 9 nm, so the quantum confinement effect is expected, and the absorption edge shifts to a higher wavelength with increasing particle size [73]. Furthermore, the insertion of  $\text{SiO}_2$ , as a shell, on the  $\text{TiO}_2$  core creates defects and causes the absorption edge to be redshifted, and the bandgap decreases accordingly. Decreasing the bandgap makes the prepared samples possible candidates in photovoltaic applications. Raising the defects and the disorder in the electronic levels of the materials could be proved by measuring the band tail width in the forbidden region through the following relation [71];

$$\alpha = \alpha_0 e^{\left[\frac{h\nu}{E_U}\right]} \quad (3)$$

where  $\alpha_0$  is constant and  $E_U$  is the band-tail width or Urbach energy. The  $E_U$  value was investigated through the linear relation between  $\ln(\alpha)$  on Y-axis and  $(h\nu)$  on X-axis (Fig.7). From Table 1, the increase in the  $E_U$  value indicates the increase in the disordering of the material and the increase of localized states in the forbidden region caused by the produced defects due to the insertion of  $\text{SiO}_2$ .

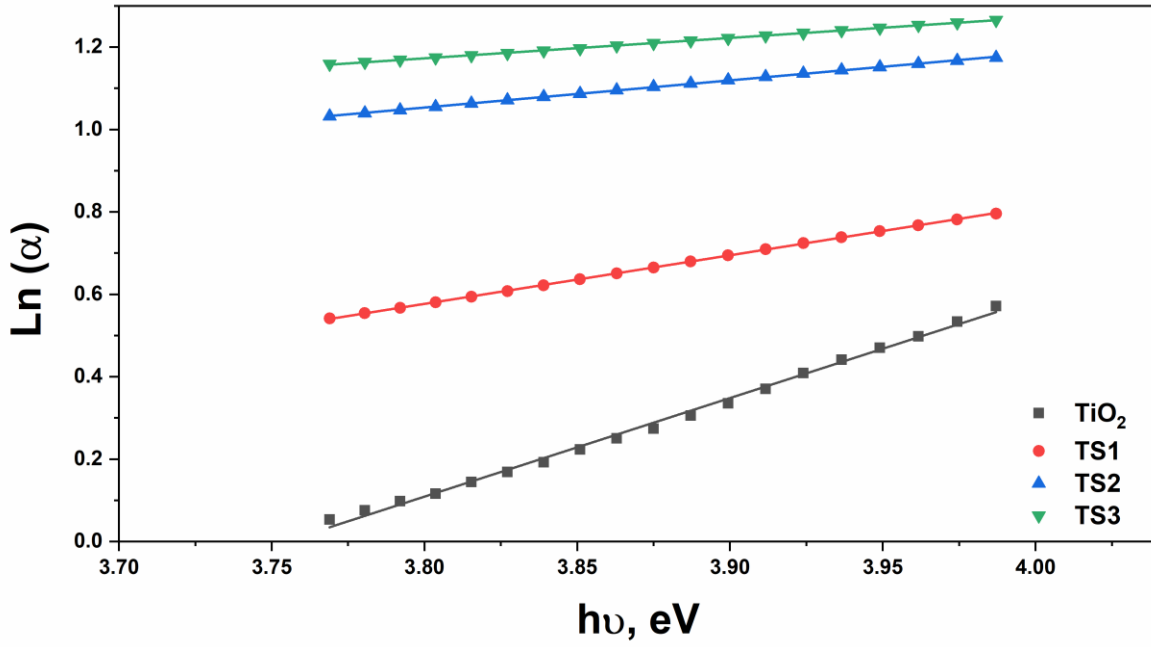


Fig.7:  $\ln(\alpha)$  against  $h\nu$  of the prepared samples.

Table 1: Optical bandgaps and Urbach energy for prepared samples.

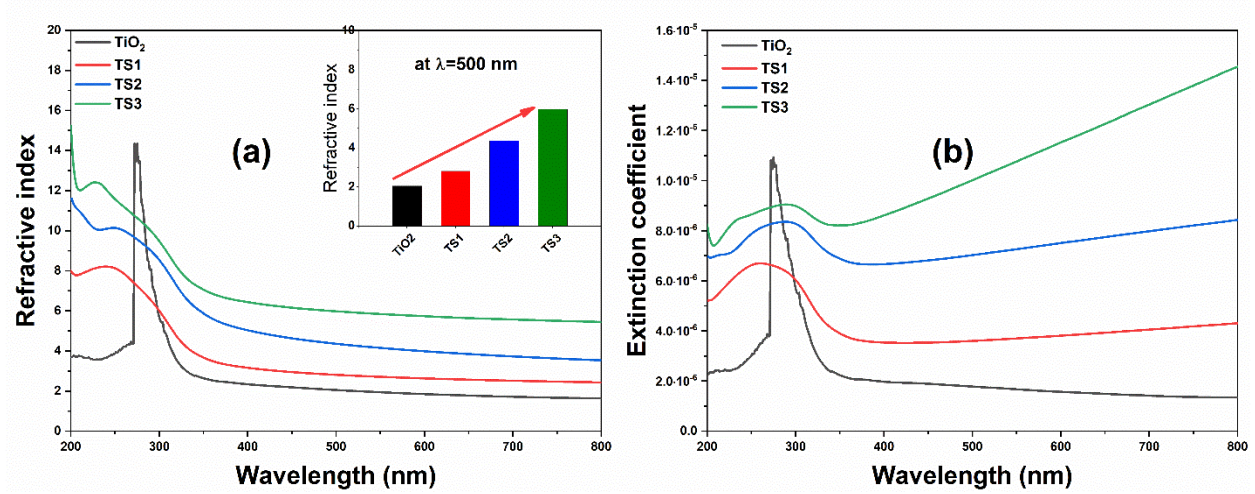
nanocomposite ID	$E_{gd}$ (eV)	$E_{gi}$ (eV)	$E_U$ (eV)	$n(\lambda=500nm)$
TiO <sub>2</sub>	3.97	3.23	0.42	2.052
TS1	3.60	2.58	0.85	2.809
TS2	3.37	1.66	1.52	4.359
TS3	3.26	1.35	2.02	5.979

The refractive index ( $n$ ) and the extinction coefficient ( $k$ ) are related to the materials' reflectivity and attenuation or absorption coefficient, respectively, and are two of the most essential optical factors. Both factors are related to each other through the following relations [71];

$$n = \frac{1-R}{1+R} + \sqrt{\frac{4R}{(1-R)^2} - k^2} \quad (4)$$

$$k = \frac{\alpha\lambda}{4\pi} \quad (5)$$

The reflectance ( $R = 1 - \sqrt{T * e^A}$ ) was investigated through the absorption (A) and transmission (T) data obtained from the UV-visible spectra. As a function of the incident light wavelength, the  $n$  and  $k$  dependencies are represented in Fig. 8.



**Fig.8: Plots of (a) Refractive index and (b) extinction coefficient as a function of wavelength for prepared samples.**

It was shown that  $n$  increases with increasing the thickness of the shell  $\text{SiO}_2$ , as illustrated in the inset of Fig.8a. Moreover, in the UV region ( $200 \leq \lambda \leq 300$  nm),  $n$  has a higher value due to the resonance effect, which is caused by the interaction of the incoming electromagnetic-radiation and the electrons polarization, as the incident photons frequency becomes equal to the plasma frequency [76, 77]. Thus, the refractive index exhibits anomalous dispersion in this region. Moving towards the visible region ( $400 \leq \lambda \leq 800$  nm), the value of  $n$  decreases exponentially. However, at higher wavelengths, the decrease in  $n$  value is slight and appears to remain constant. The absorption of the electromagnetic waves became less, and most of the incident energy was reflected. As shown in Fig.8b, based on this data, it appears that  $k$  decreases as the wavelength of incident photons increases, and seems constant for the pure  $\text{TiO}_2$  NPs sample in the visible region. The  $k$  value then begins to rise with the insertion of  $\text{SiO}_2$ , rises with the growth in shell thickness, and rises significantly at the highest shell thickness. The high value of the refractive index makes it useful as an anti-reflective coating in solar cell devices.

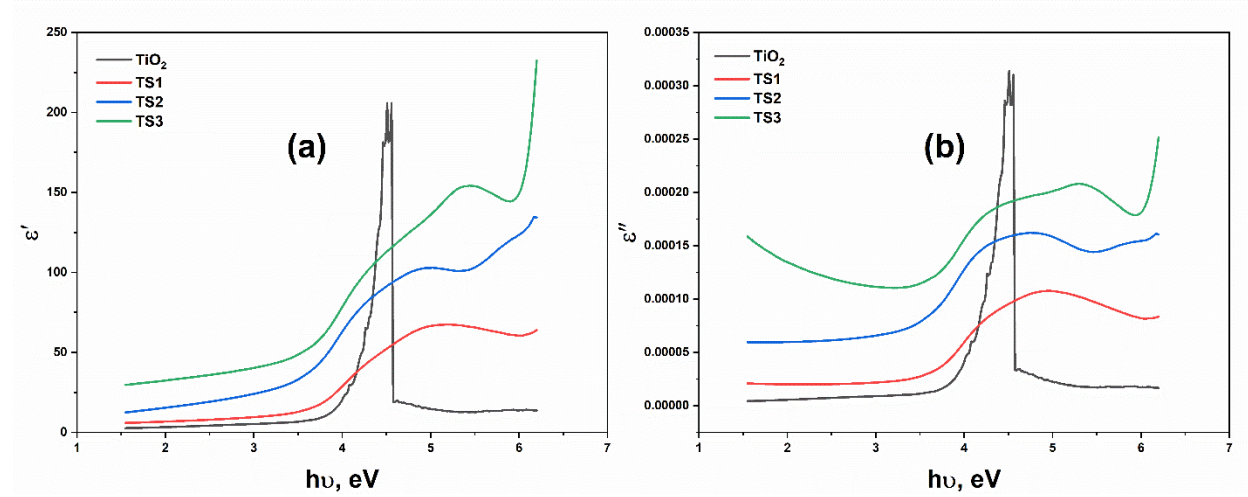
### 3.2.2. Optical dielectric constants

Optical dielectric constants, real ( $\epsilon'$ ) and imaginary ( $\epsilon''$ ) parts, depend on the  $n$  and  $k$  values of the material which can be determined through the following relations [71];

$$\epsilon' = n^2 - k^2 \quad (6)$$

$$\epsilon'' = 2nk \quad (7)$$

The dependences of both factors,  $\epsilon'$  and  $\epsilon''$ , on the incident photon energy are illustrated in nFig.9. obviously  $\epsilon'$  and  $\epsilon''$  increases with increasing the shell thickness.



**Fig.9: Plots of (a)  $\epsilon'$  and (b)  $\epsilon''$  vs.  $h\nu$  for prepared samples.**

The similarity between the behavior of the  $\epsilon'$  with refractive index and  $\epsilon''$  with  $k$  is evident. Moreover,  $\epsilon'$  value is larger than that for  $\epsilon''$ . Increasing the shell thickness causes a change in polarization, resulting in high  $\epsilon'$  and low  $\epsilon''$  values and less energy dissipation.

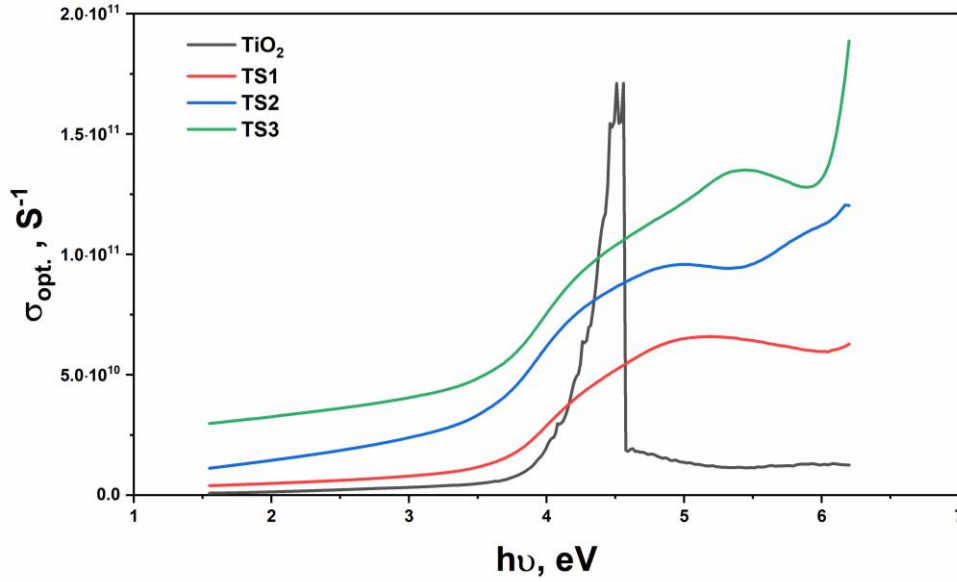
Optical conductivity ( $\sigma_{opt.}$ ) is dependent upon various factors such as  $\alpha$ ,  $n$ , and the frequency of incident photons through the following relation [77]:

$$\sigma_{opt.} = \frac{\alpha n c}{4\pi} \quad (8)$$

where ( $c=3 \times 10^8$  m/sec) is the speed of light in free space. Fig.10 depicts the correlation between the energy of incoming photons and  $\sigma_{opt.}$ . It can be observed that as the shell thickness increases the  $\sigma_{opt.}$  increases. This might be due to the increase in  $\alpha$  and  $n$  values of the material with increasing shell thickness. At high photon energies, optical conductivity is enhanced due to the excitation of electrons by incident photons. The  $\sigma_{opt.}$  value of  $\text{TiO}_2$  NPs at 600 nm is about



$1.45 \times 10^9 \text{ sec}^{-1}$  which increases to about  $5.01 \times 10^9$ ,  $1.5 \times 10^{10}$ , and  $3.31 \times 10^{10} \text{ sec}^{-1}$  for the core-shell  $\text{TiO}_2@\text{SiO}_2$  with different shell thickness TS1, TS2, and TS3, respectively.



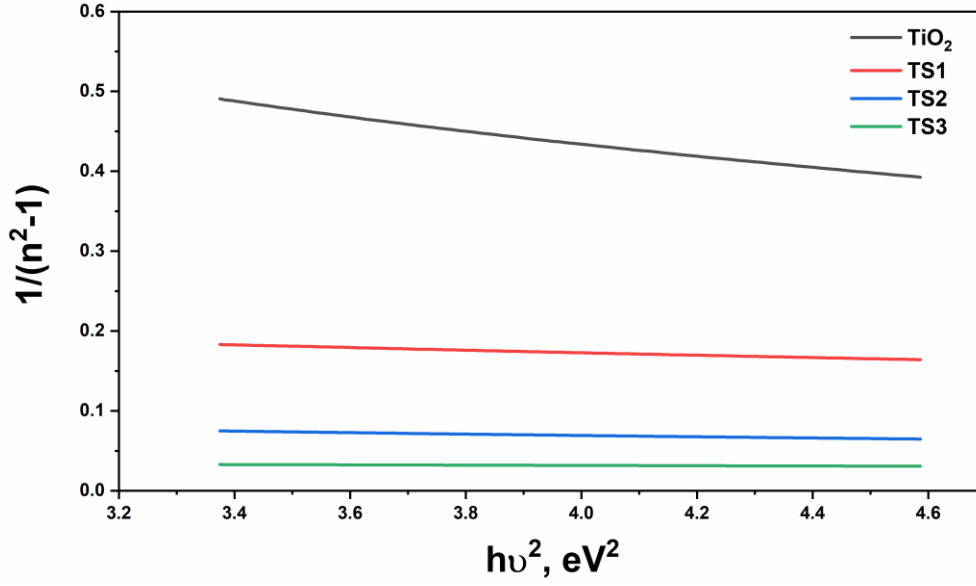
**Fig.10: The variation of  $\sigma_{\text{opt.}}$  with  $(h\nu)$  of the prepared samples.**

### 3.2.3. Dispersion energy parameters

The dispersion energy is a significant factor in researching optical materials for spectral dispersion and visual communications. The analysis of refractive index dispersion has been conducted by implementing the single oscillator concept using the Wemple-DiDomenico relationship, which depends on the dispersion-energy ( $E_d$ ) and the oscillator-energy ( $E_0$ ) as follows [78, 68]:

$$\frac{1}{(n^2-1)} = \frac{E_0}{E_d} - \left(\frac{1}{E_0 E_d}\right)(h\nu)^2 \quad (9)$$

Plot of  $1/(n^2 - 1)$  versus  $(h\nu)^2$  is illustrated in Fig. 11. From the slope and the intercept values, we can derive the values of  $E_d$  and  $E_0$  (Table 2).



**Fig.11:**  $1/(n^2-1)$  vs.  $(h\nu)^2$  for the prepared samples.

**Table 2:** The optical dispersion parameters of the  $\text{TiO}_2$  and  $\text{TiO}_2@\text{SiO}_2$  core-shell.

Sample ID	$E_d(\text{eV})$	$E_0(\text{eV})$	$f, \text{eV}^2$	$n_0$	$\epsilon_s = n_0^2$	$\chi^{(1)} (\text{esu})$	$\chi^{(3)} (\text{esu})$	$n_2 (\text{esu})$
<b>TiO<sub>2</sub></b>	4.06	3.07	12.45	1.523	2.32	0.11	$2.08 \times 10^{-14}$	$5.14 \times 10^{-13}$
<b>TS1</b>	16.40	3.87	63.49	2.289	5.24	0.34	$2.20 \times 10^{-12}$	$3.62 \times 10^{-11}$
<b>TS2</b>	34.02	3.50	119.19	3.272	10.71	0.77	$6.06 \times 10^{-11}$	$6.97 \times 10^{-10}$
<b>TS3</b>	122.58	4.74	581.40	5.181	26.85	2.06	$3.04 \times 10^{-09}$	$2.21 \times 10^{-08}$

As shown in Table 2, as the shell thickness increases around the core, the  $E_0$  and  $E_d$  values increase. These values can be used to determine  $n_0$  (static-refractive index),  $\epsilon_s$  (static-dielectric constant), and  $f$  (optical-oscillator strengths) for optical transitions, which are calculated according to the following relations [68];

$$\text{at } h\nu = 0 \rightarrow n_0 = \sqrt{1 + \frac{E_d}{E_0}} \quad (10)$$

$$\epsilon_s = n_0^2 \quad (11)$$

$$f = E_d E_0 \quad (12)$$

The calculated values of  $n_0$ ,  $\epsilon_s$ , and  $f$  are extracted in Table 2. As the thickness of the shell increases,  $n_0$ ,  $\epsilon_s$ , and  $f$  values increase.

The nonlinear parameters like; first- ( $\chi^{(1)}$ ), third- ( $\chi^{(3)}$ ) order nonlinear optical susceptibility, and nonlinear refractive index ( $n_2$ ) were evaluated as follows [79];

$$\chi^{(1)} = \frac{E_d}{4\pi E_o} \quad (13)$$

$$\chi^{(3)} = 6.82 \times 10^{-15} \left[ \frac{E_d}{E_o} \right]^4 \quad (14)$$

$$n_2 = \frac{12\pi\chi^{(3)}}{n_o} \quad (15)$$

Table 2 summarizes the calculated values. From Table 2, as the shell thickness increases, the  $\chi^{(1)}$ ,  $\chi^{(3)}$ , and  $n_2$  values increase. The  $\chi^{(3)}$  value for pure TiO<sub>2</sub> NPs is about  $2.08 \times 10^{-14}$  increases to about  $3.04 \times 10^{-09}$  for core-shell TS3 sample. This value is better than that represented for TiO<sub>2</sub>/ZnO thin film ( $3.97 \times 10^{-12}$ ) [80]. The core-shell TiO<sub>2</sub>@SiO<sub>2</sub> structure is the perfect choice for photonic and nonlinear optical devices due to its improved nonlinear optical parameter. A high nonlinear refractive index can be used in communication fibers and optical limiting devices.

The  $n_2$  value increases as the  $n$  value increases and  $E_g$  decreases, directly correlating with an increase in metallicity. A correlation between the energy gap of oxides and their molar refraction was suggested by Duffy and Lorentz–Lorentz equations as follows [80-82];

$$E_g = 20 \left( 1 - \frac{R_m}{V_m} \right)^2 \quad (16)$$

$$1 - \frac{R_m}{V_m} = \left( \frac{E_g}{20} \right)^{\frac{1}{2}} = 1 - \frac{n_o^2 - 1}{n_o^2 + 2} \quad (17)$$

where  $R_m$  and  $V_m$  are the molar refraction and volume, respectively. The value of metallization criterion ( $1 - R_m/V_m$ ) was calculated based on  $E_g$  and  $n_o$  values and summarized in Table 3.

**Table 3:** Metallization criterion values.

Sample ID	$\left( \frac{E_g}{20} \right)^{\frac{1}{2}}$	$1 - \frac{n_o^2 - 1}{n_o^2 + 2}$
<b>TiO<sub>2</sub></b>	0.45	0.69
<b>TS1</b>	0.42	0.41
<b>TS2</b>	0.41	0.24
<b>TS3</b>	0.40	0.10

The metallization criterion of TiO<sub>2</sub> NPs is about 0.45, which is near the value in the literature (0.39) [81]. This is due to the difference in bandgap associated with the tiny size of the as-prepared TiO<sub>2</sub>. The metallization criterion value decreases with increasing the shell thickness. The width of both valence and conduction bands becomes large when the  $R_m/V_m$  value is high, resulting in a narrow band gap and increasing the metallicity of the material [80, 81]. These values provide a solid foundation for anticipating novel nonlinear optical materials.

#### **4. Conclusion**

TiO<sub>2</sub> and core-shell TiO<sub>2</sub>@SiO<sub>2</sub> were synthesized utilizing sol-gel and Stöber methods. The shell thickness was controlled by controlling TEOS concentration during the reaction process. XRD showed that the crystallinity of TiO<sub>2</sub> decreased upon encapsulation with an amorphous silica shell and became more amorphous as the shell thickness increased. TEM analysis confirms the formation of the core-shell TiO<sub>2</sub>@SiO<sub>2</sub> with different shell thicknesses. TiO<sub>2</sub> nanoparticle absorption increases by encapsulation with a silica shell and increases as the shell thickness increases. The absorbance of the TiO<sub>2</sub>@SiO<sub>2</sub> nanocomposite can be adjusted by varying the silica shell thickness. It is imperative to adjust this parameter for optimum results. The investigation showed that as the shell thickness increases, the refractive index, optical dielectric constants, and optical conductivity increase, while the optical bandgap and metallization criterion decrease. The core-shell TiO<sub>2</sub>@SiO<sub>2</sub> structure is the perfect choice for photonic and nonlinear optical devices due to its improved nonlinear optical parameter.

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**Declaration of interests**

☒ The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

☐The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: